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The integration of (A 3) is straightforward, although not as simple as before. The resulting change of wave in wave amplitude over the slope is

$$\frac{\eta}{\eta_0} = A[h(x)]\exp\left\{\frac{1}{s}\left(\frac{\nu}{2\omega}\right)^{\frac{1}{2}}\left[\left(\frac{\omega^2}{s}+\frac{2}{b}\right)(kh-k_0h_0)-2(k-k_0)\right]\right\}, \qquad \text{(A 4)}$$

which is identical with the previous result except for the factor 2 before the last term on the right.

Discussion

To appraise the change accorded to these corrections, the machine computation of (A 4) was rerun for the cases where the largest differences would occur; i.e. for the smallest experimental values of kh ($h = 1$ cm and $0{\cdot}1 < k < 0{\cdot}5$) with the result that the maximum difference encountered was about 10 %. At the 5 cm depth the maximum error was 3 %. The percentage error in the previous formulation, of course, becomes larger the lower the frequency, but the entire dissipation becomes negligible at low frequency. Therefore the principal results of the experiment are still valid.

2. 'A uniformly valid solution for the hypersonic flow past blunted bodies', by W. Schneider, *J. Fluid Mech.* vol. **31**, 1968, p. 397.

(i) Equation (53*b*): The characters immediately before and immediately behind the colons should be exchanged.

(ii) Figure 5: The words (shock) and (body) should be exchanged.

(iii) Page 406, line 24: (32) should read (39).

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